Transition métal-non métal dans le composé $(BEDT-TTF)_8Hg_4Br_{12}(C_6H_5Br)_2$

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Motivation: (BEDT-TTF)₈Hg₄X₁₂(C₆H₅Y)₂ family, where (X, Y) = (Cl, Br)



X = CI: metallic ground-state X = Br: metal-non metal transition even though the four members of this family are isostructural at room temperature

R. Lyubovskaya et al. Synth. Met. 55-57 2899 (1993)

 $(BEDT-TTF)_{8}Hg_{4}X_{12}(C_{6}H_{5}Y)_{2}$: crystalline structure



X = Y = Br (150 K)

⁽conducting planes: (BEDT-TTF)⁰⁵⁺ cation layers



Insulating planes: HgX₃⁻ anion + C₆H₅Y (solvent)

$(BEDT-TTF)_{8}Hg_{4}X_{12}(C_{6}H_{5}Y)_{2}$: crystalline structure



 $(BEDT-TTF)_8Hg_4X_{12}(C_6H_5Y)_2$: crystalline structure



Electronic structure (Extended Hückel)



Two pairs of crossing q-1D sheets

Degeneracy removals yields compensated electron and hole orbits

→2D Fermi surface

X = Cl, Y = Br: metallic ground-state

X = Y = Br: metal non-metal transition at ~150 K. Hidden nesting → density wave ?







Quantum oscillations closed orbits --> test for Fermi surface

- Periodic in 1/B
- Frequency proportional to orbit area (in k-space)
- Temperature and field-dependence yield physical parameters

Interlayer magnetoresistance in pulsed magnetic field (55 T) → Shubnikov-de Haas oscillations





Applied pressure powerful tool for tuning physical properties

Need for isothermal transport measurements at liquid helium temperatures in pulsed fields

Pressure cell: zirconia anvils – Pmax = 1.2 GPa gasket anvil track (a) compression (b) chamber 10 mm Gasket's creep during the load phase 18nn

M. Nardone et al. Cryogenics 41 175 (2001)

Ambient pressure data for X = CI: metallic ground state

Fourier analysis



a lot of frequencies are observed

Vignolles et al. Eur. Phys. J. B **31** 53 (2003)

Ambient pressure data for X=CI, Y=Br (metallic ground state)

Linear combinations of the frequencies linked to the basic electron and hole orbits (a) and of the δ and Δ pieces .



Model system for the study of quantum oscillations in networks of coupled orbits

Vignolles et al. Eur. Phys. J. B **31** 53 (2003)

Derivation of physical parameters: Lifshits-Kosevich formalism

Fourier amplitude: A $\propto R_T R_D R_{MB}$

at a given field value: Fermi-Dirac smearing $A \propto R_T = \frac{Tm^*}{B} / sinh \left(\frac{u_0 Tm^*}{B}\right)$

temperature dependence yields effective mass (m*)



Derivation of physical parameters: Lifshits-Kosevich formalism

Fourier amplitude: A $\propto R_T R_D R_{MB}$

At a given temperature:

Dingle damping factor: $R_D = \exp\left(-\frac{u_0 T_D m^*}{B}\right)$ $T_D = \frac{\hbar}{2\pi k_D \tau}$ τ : scattering time Magnetic breakdown damping factor: $R_{MB} = [1 - exp(-B_0/B)]^2$ B_0 : MB field 0.03 T = 1.65 KFourier amplitude $B_0 = (35 \pm 10) T$ $T_{D} = (0.2 \pm 0.2) \text{ K}$ $\tau^{-1} = (1.7 \pm 1.7) \times 10^{11} s^{-1}$ very good crystals! 0.00 10 20 30 40 50 0 B (T)

Effect of applied pressure: zero-field resistance



X = Br: pressure-induced metallic state





X = Cl: bump under pressure Defects and (or) correlations X = Br at low pressure: hysteresis at high pressure: metallic state

Effect of applied pressure: zero-field resistance at low temperature



X = Cl and X = Br at high pressure: $R = R_0 + AT^2$ \rightarrow correlated Fermi liquid A decreases as applied pressure increases Effect of applied pressure: SdH oscillations

X = Cl (metallic ground state)







Still many frequency combinations (magnetic breakdown) Same features as at ambient pressure Effect of applied pressure: SdH oscillations spectra

X = Br (pressure-induced metallic ground state)





Almost no frequency combinations: suggests magnetic breakdown gap larger than for X = Cl

Effect of applied pressure: Oscillations frequency linked to basic orbits



- organic metals
- Orbits area in good agreement with band structure calculations

9.6 % of the FBZ area

Effect of applied pressure: Magnetic breakdown field







Magnetic breakdown not observed for X = Br Accounts for the absence of frequency combinations. Consistent with band structure calculations

Effect of applied pressure:



Strong decrease of the MB field

Strong (and reversible) increase of the scattering rate

may explain zero-field resistance:



Effect of applied pressure: Effective mass and zero-field resistance



Strong decrease of both the effective mass and the coefficient (A) of the T² law of the zero-field resistance.

Effect of applied pressure: Effective mass and zero-field resistance

Low temperature resistivity driven by electron correlations: $\mathbf{A} \propto \mathbf{m}^{*2}$



In the framework of the Brinkman-Rice scenario a divergence of the effective mass as approaching a Mott transition should be observed

Mott insulator state already reported for phase diagrams of organic conductors, although mainly in κ -(BEDT-TTF)₂X salts, with different Fermi surface



Kagawa et al. Nature 436 534 (2005)

Caulfield et al. J.P.Cond.Mat **6** 2911 (1994)

Summary

Good agreement with band structure calculations in the metallic state (Extended Hückel) Zero-field interlayer resistivity:

- typical of strongly correlated Fermi liquid at low temperature (T² law)
- contribution of disorder under applied pressure at higher temperatures
- Effective mass strongly decreases as pressure increases:

could suggest a Brinkman-Rice scenario (divergence of m*) for X =Br

Next steps

For X = Br:

X-ray data at low temperature (charge density wave, charge ordering?) Magnetoresistance data in the pressure range 0.3 - 0.7 GPa (i.e. within the metal-non metal transition) to check the behaviour of the effective mass